

Soft Matter

Lecture 6



Today's lecture (6): Mechanical properties of soft matter

- Deformation: elongation and shear
- Strain, strain rate and stress
- Elastic solids and viscous liquids and viscoelasticity
 - Viscosity, Maxwell model, Kelvin-Voigt model
- Rheology: instrumentation and measurement types
- Strange fluids: some soft matter examples
 - Shear thinning, yield stress, shear thickening, thixotropy

Deforming soft matter in daily life

painting



mixing

squirting



squeezing

The blood miracle of Saint Gennaro (Januarius)

- Annual liquefaction of blood kept in sealed glass ampoule, Naples



2005



2013



2019

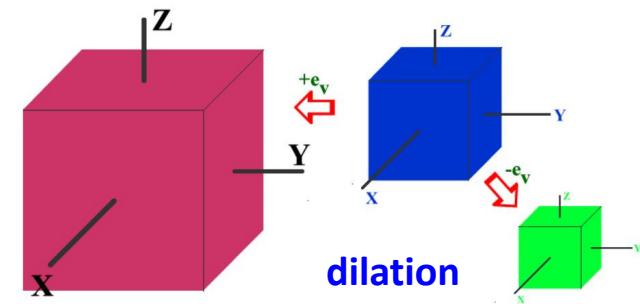
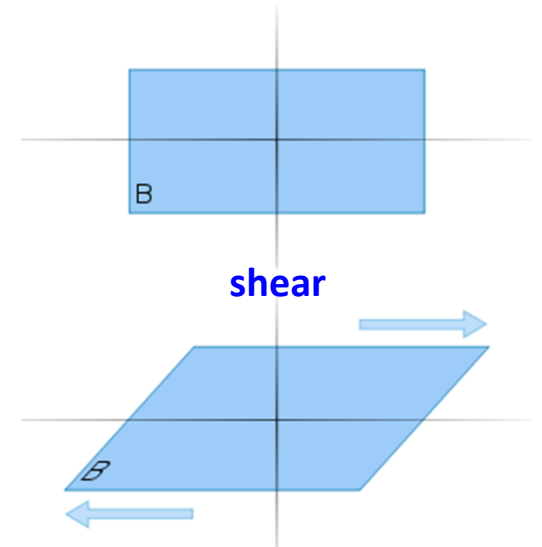
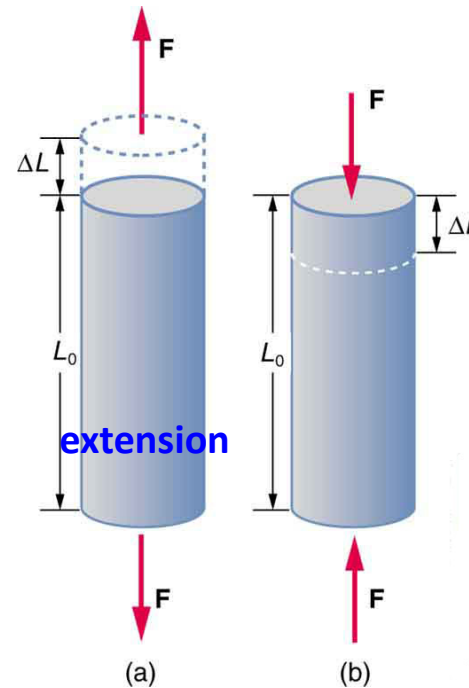


Das Blut des heiligen Januarius.
Nach der Natur gezeichnet von C. Grob.
(Hälfte der natürlichen Größe.)

Deformation

- Change in the metric properties of a continuous body (e.g., a material)
- Useful to elucidate material behavior
 - Linear/extension
 - Plane
 - Volume

 - Uniform extension
 - Simple shear
 - Pure dilation/compression (overall shape remains the same)



Strain and stress

- **Strain**: deformation per unit length (dimensionless)

$$\gamma = \Delta x / y$$

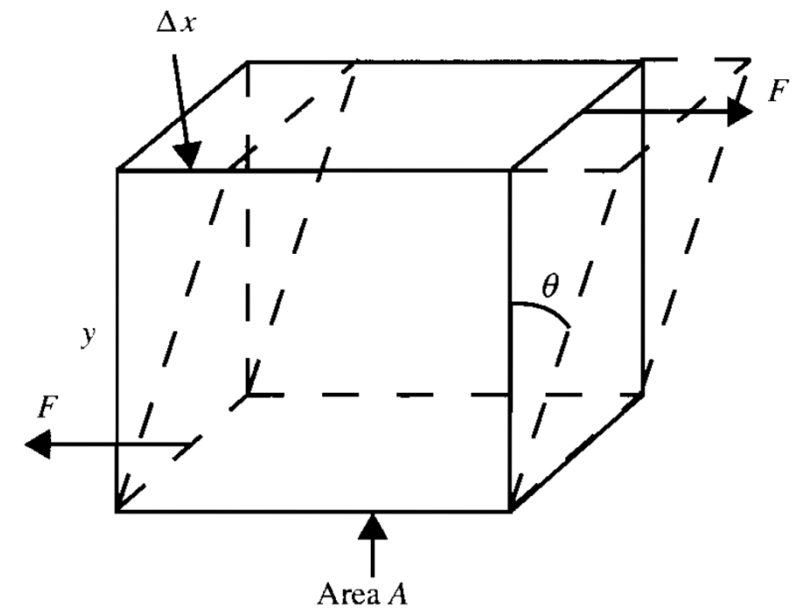
- Time derivative of strain = **strain rate** (s^{-1})

$$\dot{\gamma} = d\gamma / dt = v_x / y$$

- **Stress**: force per unit area ($N/m^2 = Pa$)

$$\sigma = F / A$$

- Depending on direction, one can have:
 - Elongational strain / elongational stress
 - Shear strain / shear stress
 - Shear (strain) rate



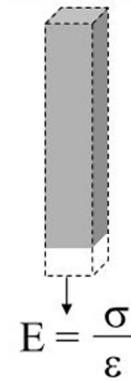
Simple materials

- Hookean solids: elastic deformation

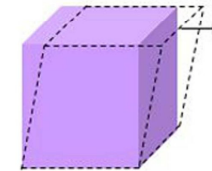
$$\sigma = G \gamma$$

- Modulus G (Pa)

Young's Modulus

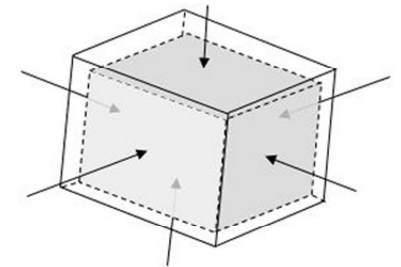


Shear Modulus



$$G = \frac{\tau}{\gamma}$$

Bulk Modulus



$$B = \frac{\sigma_{\text{hyd}}}{\Delta V/V_0}$$

- Newtonian liquids: plastic deformation (viscous)

$$\sigma = \eta \dot{\gamma}$$

- Viscosity η (Pa·s)

Soft matter is really “soft”

- Hookean solids: elastic deformation

$$\sigma = G \gamma$$

– Modulus G (Pa) = energy per unit volume ($\text{N m}^{-2} = \text{N} \cdot \text{m} \cdot \text{m}^{-3} = \text{J m}^{-3}$)

– Crystal $E = 1 \text{ eV}$ (96 kJ/mol), $l = 0.1 \text{ nm}$, $G = 100 \text{ GPa}$

– Soft matter $E = 1 k_B T$ (2.4 kJ/mol), $l = 100 \text{ nm}$, $G = 10 \text{ Pa}$



10^{-10}

Viscosity and soft matter characteristics

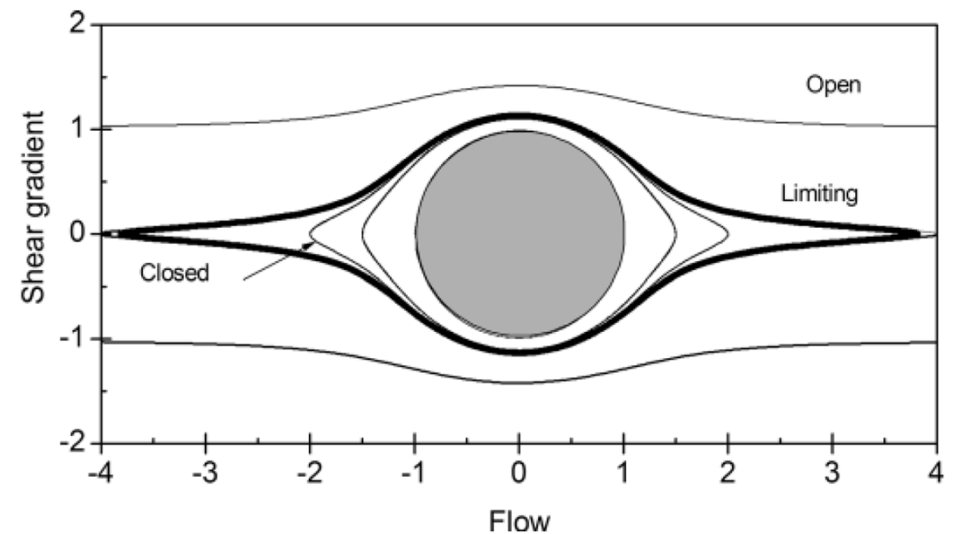
- Newtonian liquids: plastic deformation (viscous)

$$\sigma = \eta \dot{\gamma}$$

- Dispersed particles increase the viscosity
- Hydrodynamic effect

- Einstein's equation for suspension rheology:

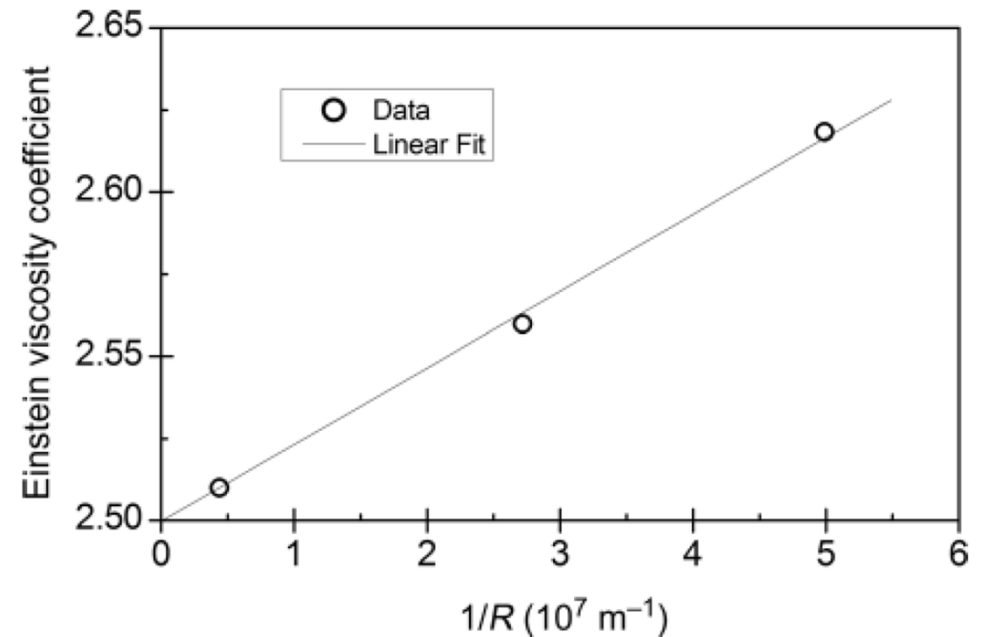
$$\eta = \eta_0(1 + 2.5 \phi)$$



Viscosity and soft matter characteristics

- Einstein's equation for suspension rheology

$$\eta = \eta_0(1 + 2.5\phi)$$



Brodnyan 1968

Viscosity of semidilute suspensions

- Einstein's equation for suspension rheology

$$\eta = \eta_0(1 + 2.5\phi)$$

- are the first two term from a Taylor series for the suspension viscosity:

$$\eta = \eta_0(1 + 2.5\phi + c_2\phi^2 + c_3\phi^3 + \dots)$$

- Batchelor and Green (1972) showed that for hard spheres, three terms is accurate up to crystallization:

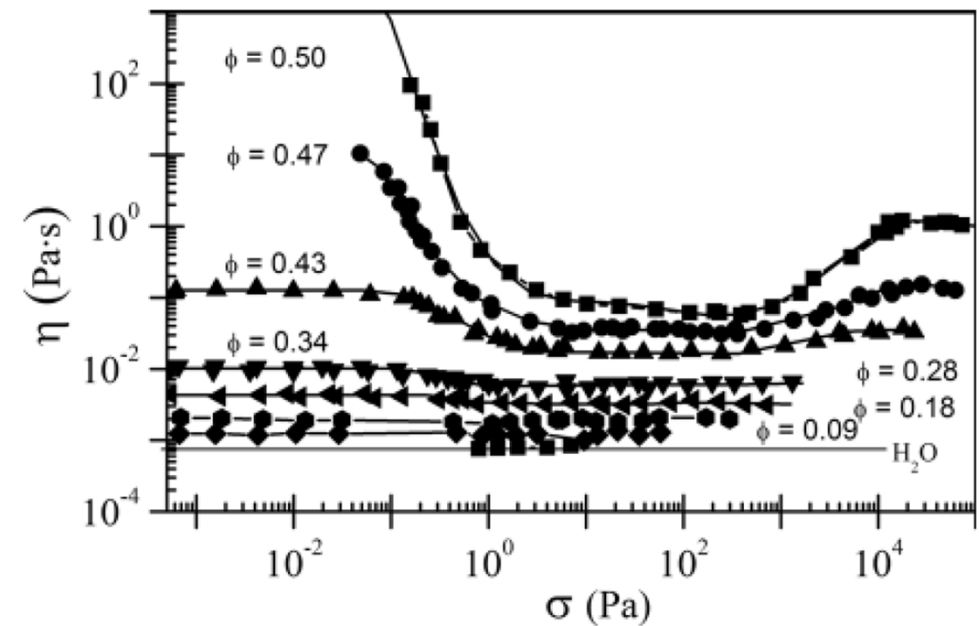
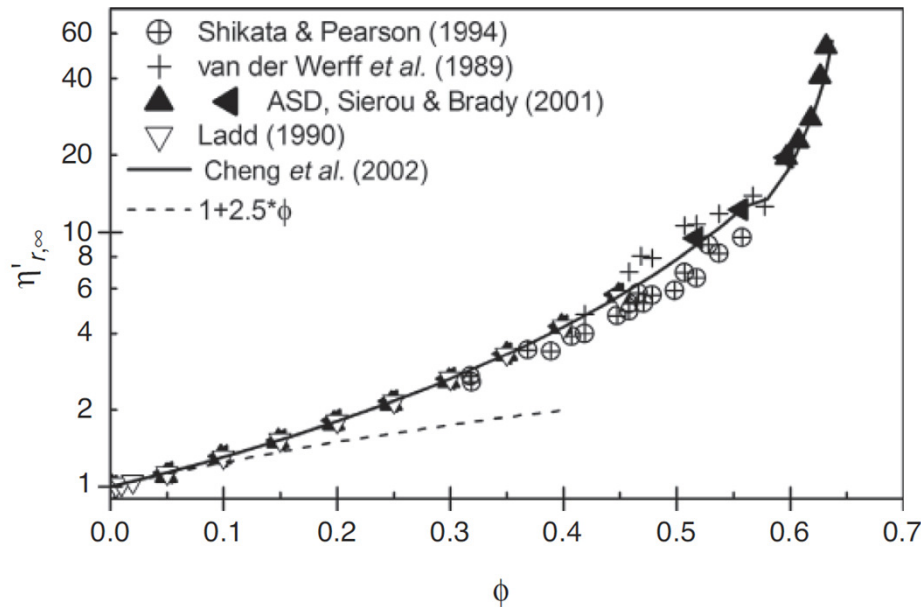
$$\eta = \eta_0(1 + 2.5\phi + 7.6\phi^2)$$

Viscosity of concentrated suspensions

- Increasingly difficult to flow

$$\eta'_{r,\infty} = \frac{1 + \frac{3}{2}\phi [1 + \phi (1 + \phi - 2.3\phi^2)]}{1 - \phi [1 + \phi (1 + \phi - 2.3\phi^2)]}, \quad 0 \leq \phi \leq 0.56,$$

$$= 15.78 \ln \left(\frac{1}{1 - 1.160\phi^{1/3}} \right) - 42.47, \quad 0.60 \leq \phi \leq 0.64.$$



On viscosity

- Viscosity η_s $\eta_0(1 + 2.5\phi)$
- Relative viscosity $\eta_r = \eta_s/\eta_0$ $1 + 2.5\phi$
- Specific viscosity $\eta_{sp} = \eta_r - 1 = \frac{\eta_s - \eta_0}{\eta_0}$ 2.5ϕ
- Reduced viscosity $\eta_{red} = \frac{\eta_{sp}}{C}$ $2.5/\rho$
- Intrinsic viscosity $[\eta] = \lim_{C \rightarrow 0} \frac{\eta_{sp}}{C}$ \Rightarrow Tutorial 6

Viscosity of dilute polymer solutions

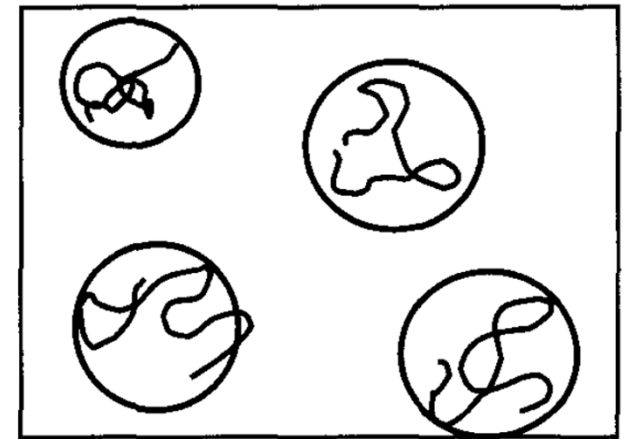
- Viscosity $\eta_s = \eta_0(1 + 2.5\phi)$

- Hydrodynamic size of polymer coil:

$$\eta_{sp} = 2.5\phi = 2.5 \left(\frac{n_p}{V} \right) V_h$$

$$\eta_{sp} = 2.5 \left(\frac{4\pi N_A C}{3M_w} \right) R_h^3 \sim R_h^3$$

$$[\eta] = \frac{2.5 N_A V_h}{M_w}$$



Dilute ($\phi < \phi^*$)

ϕ = volume fraction of particles/polymers

n_p = number of particles/polymer coils

V = volume

V_h = hydrodynamic volume of a
particle/polymer coil

R_h = hydrodynamic radius

C = mass concentration (kg/m^3)

N_A = Avogadro's number

Viscosity of dilute polymer solutions

- Hydrodynamic size of polymer coil:

$$[\eta] = \frac{2.5N_A V_h}{M_w}$$

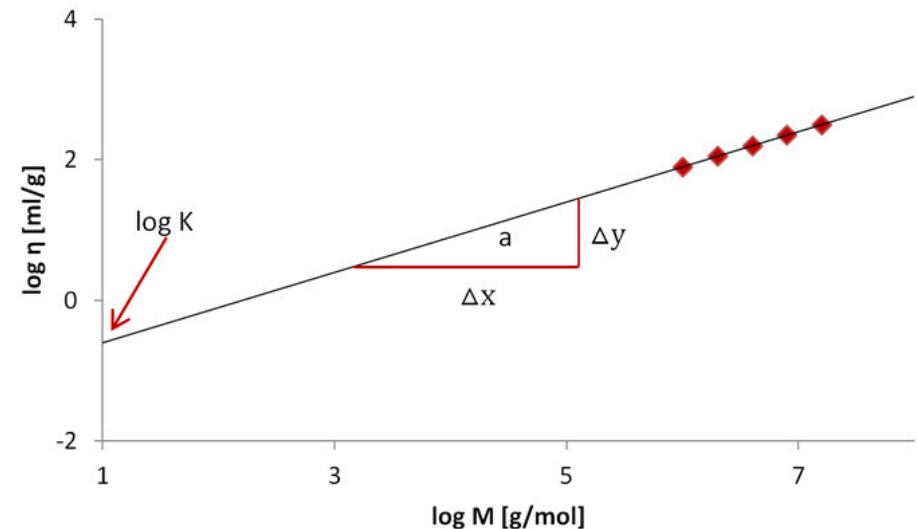
- Combine with polymer size scaling:

$$R \sim N^\nu \sim (M_w)^\nu$$

- Mark-Houwink-Sakurada relationship:

$$[\eta] = K_M (M_w)^a$$

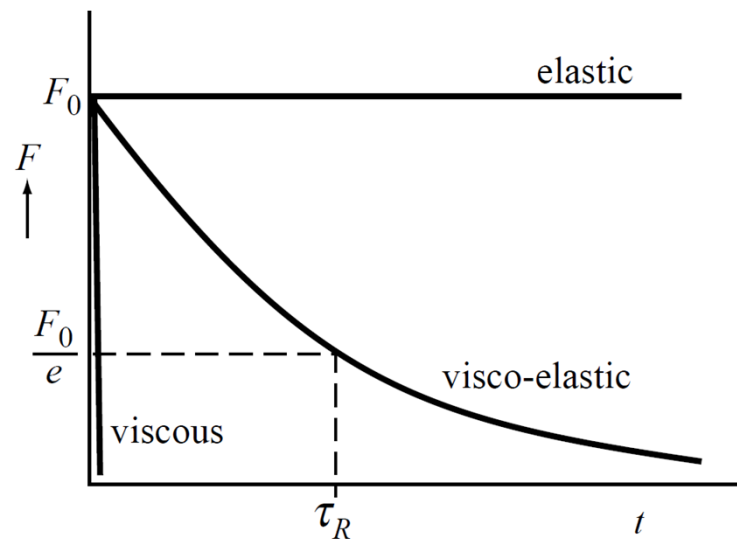
=> Tutorial 6



Ideal chain:	$a = 0.5$
Real chain in a good solvent	$a = 0.76$
Flory approximation:	$a = 0.8$

Viscoelasticity

- Response to an applied stress or strain in a time-dependent manner
- Combination of viscous and elastic response
- Typically: elastic at short timescales, viscous at long timescales, relaxation time τ defines crossover
- *Panta rhei*: everything flows



Viscoelasticity

- Maxwell model: a viscoelastic liquid
 - Spring and dashpot in series

$$F = k_e x_e = k_v \left(\frac{dx_v}{dt} \right)$$

$$x_t = x_e + x_v$$

$$\frac{x_t - x_v}{x_t} = \exp\left(-\frac{k_e}{k_v} t\right)$$

$$x_v = 0 \text{ (at } t = 0)$$

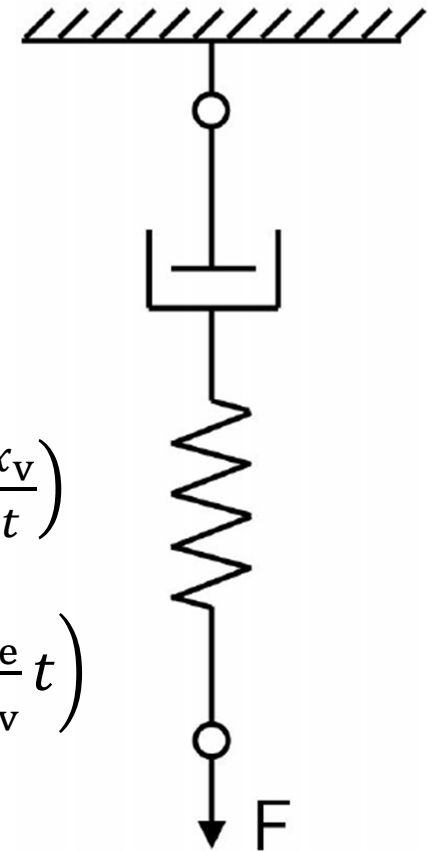
$$k_e(x_t - x_v) = k_v \left(\frac{dx_v}{dt} \right)$$

$$F(t) = k_e x_t \exp\left(-\frac{k_e}{k_v} t\right)$$

Converting to stress and strain gives:

$$\sigma(t) = G_0 \gamma \exp(-t/\tau)$$

with: $\eta_0 = G_0 \tau$



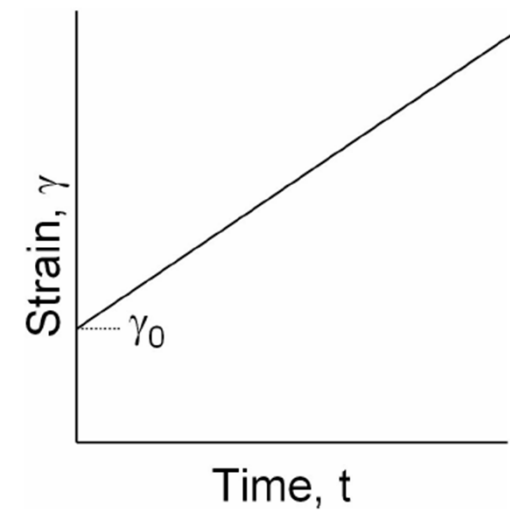
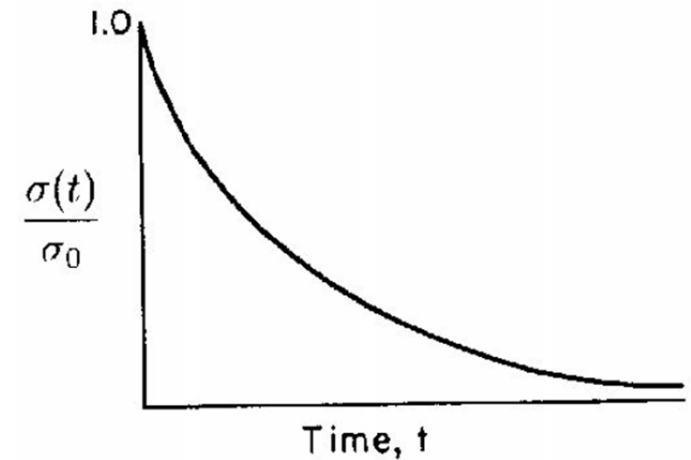
Viscoelasticity

- Maxwell model: a viscoelastic liquid

- Spring and dashpot in series

- Stress relaxation: $\sigma(t) = G_0 \gamma \exp(-t/\tau)$

- Creep: $\gamma(t) = \gamma_0(1 + t/\tau)$



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Viscoelasticity

- Kelvin-Voigt model: a viscoelastic solid
 - Spring and a dashpot in parallel

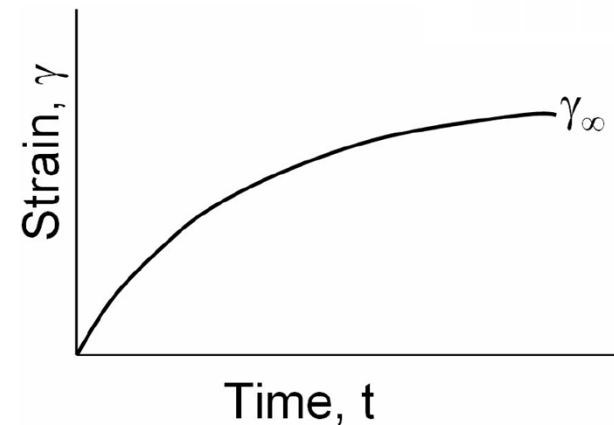
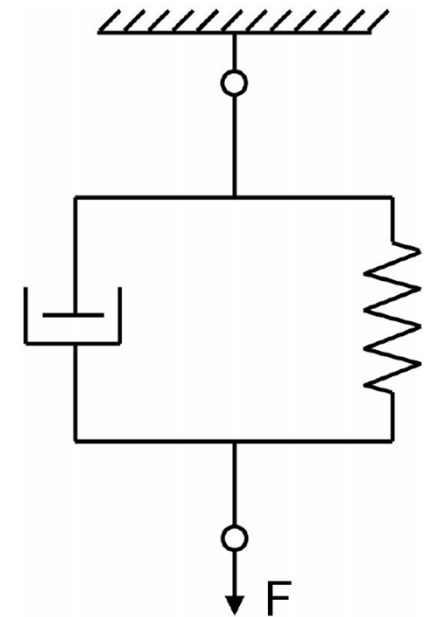
$$F = k_e x + k_v \left(\frac{dx}{dt} \right)$$

$$x = 0 \text{ (at } t = 0 \text{)}$$

$$\frac{k_e}{k_v} dt = \frac{1}{F/k_e - x} dx$$

$$x(t) = \left(\frac{F}{k_e} \right) \left(1 - \exp \left(-\frac{k_e}{k_v} t \right) \right)$$

$$\gamma(t) = \gamma_\infty (1 - \exp(-t/\tau))$$



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Fantastic materials and where to find them

- Many soft materials are viscoelastic, and show non-linear response
- Response depends on stress, rate of deformation, time, and the history of previous deformations....
- A recipe for “strange” behaviors



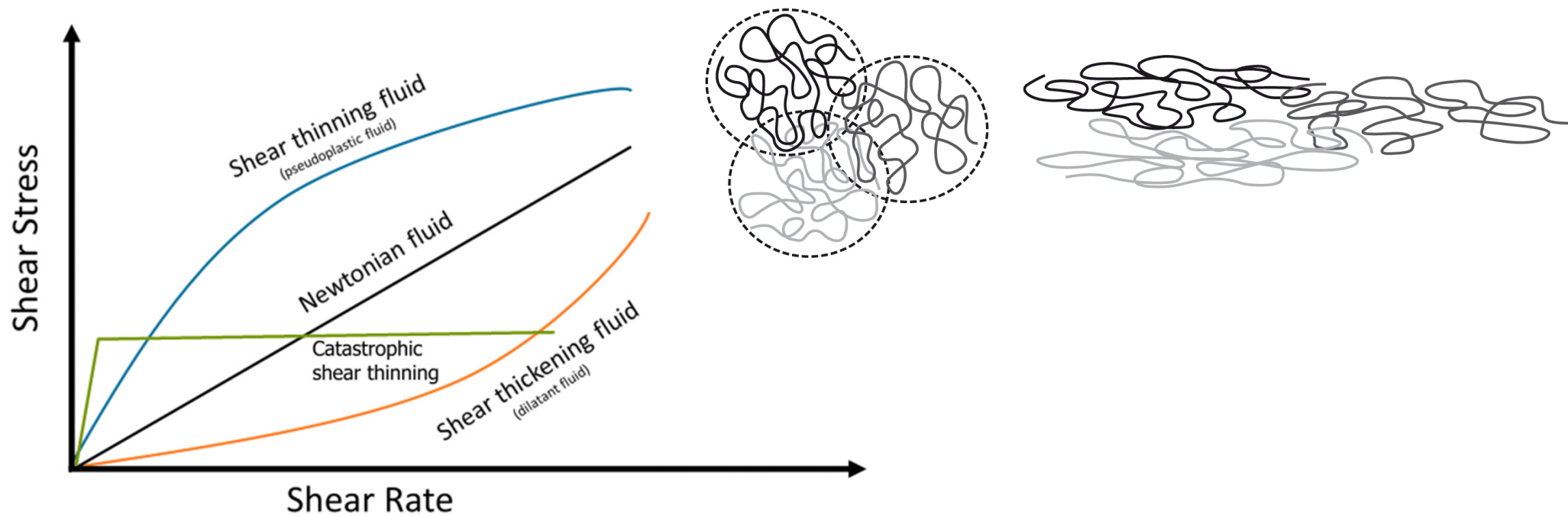
Fantastic materials and where to find them

- Shear thinning (plasticity)
- Shear thickening (dilatancy)
- Strain stiffening
- Yield stress (Bingham fluid)
 - No deformation below a threshold stress

- Thixotropy: time-dependent decrease in viscosity
 - Fluidize when shaken
- Rheopexy: time-dependent increase in viscosity
 - Solidify when shaken

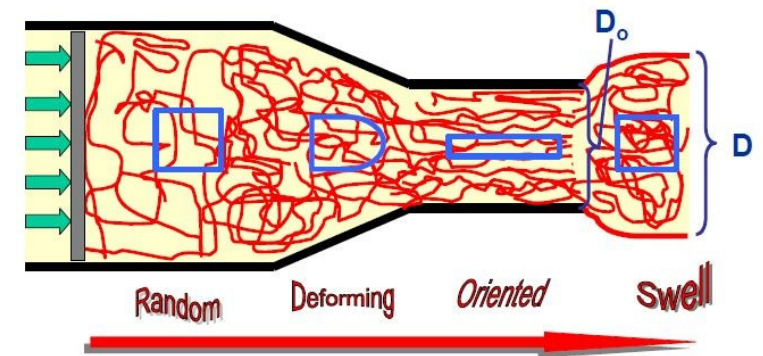
Shear thinning

- Common property of polymer solutions



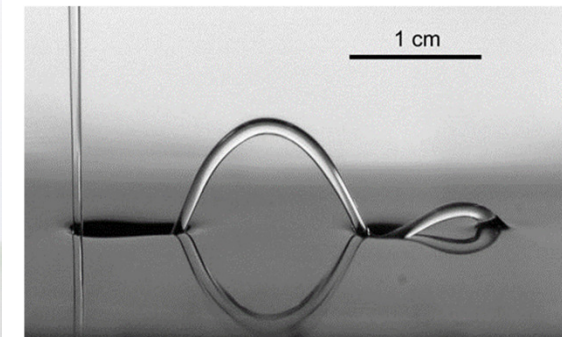
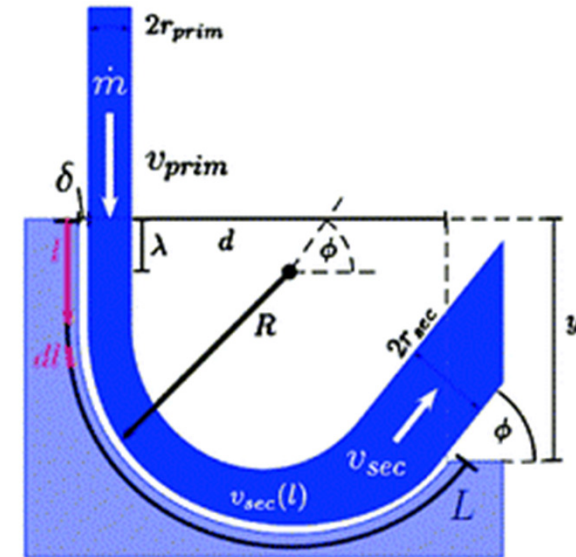
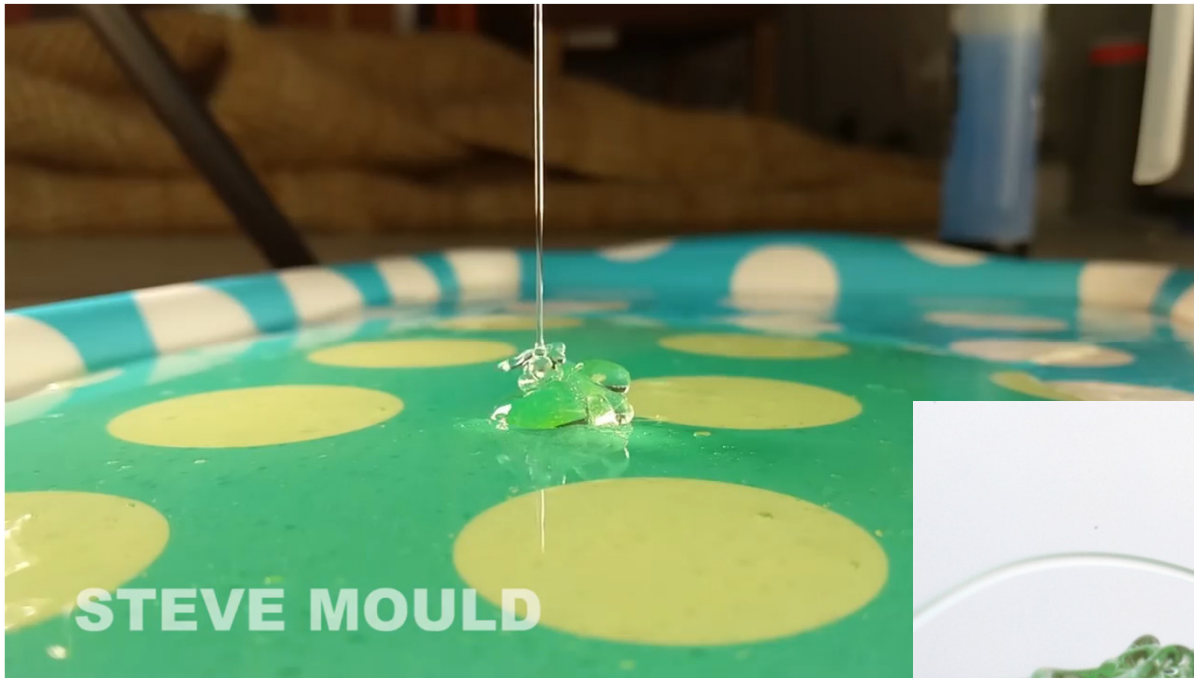
Shear thinning: Barus effect

- Barus effect (die swelling) of compressed polymer solution



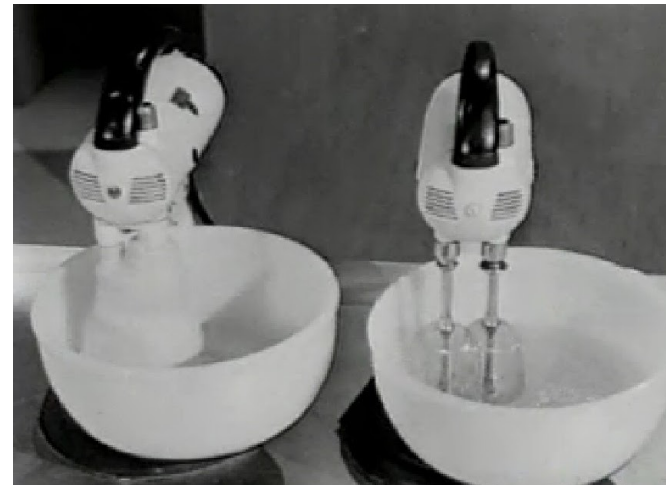
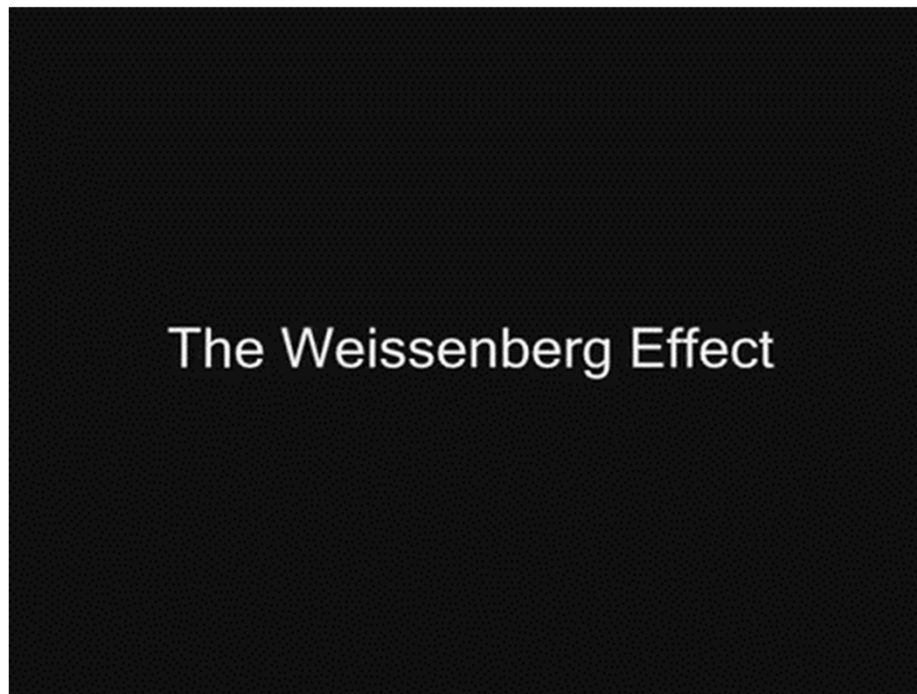
Shear thinning: Kaye effect

- Shear thinning on a tilted surface



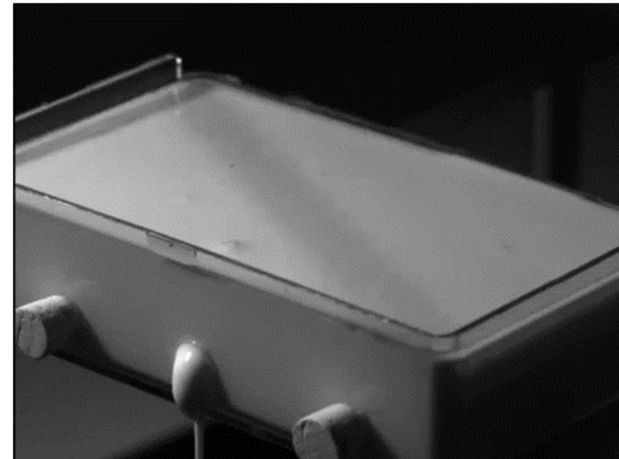
Shear thinning: Weissenberg effect

- Normal forces due to shear



Shear thickening

- Can be observed in suspensions of particles (spherical or nonspherical)

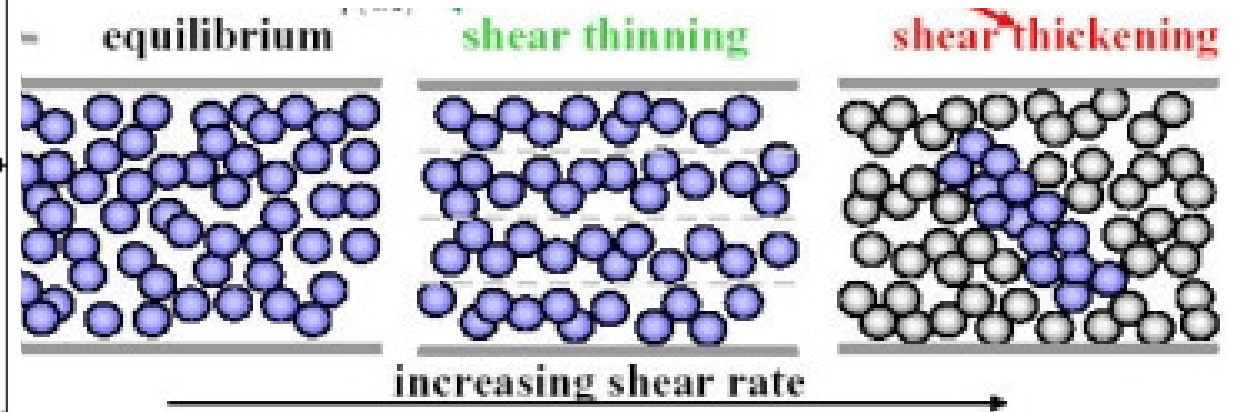
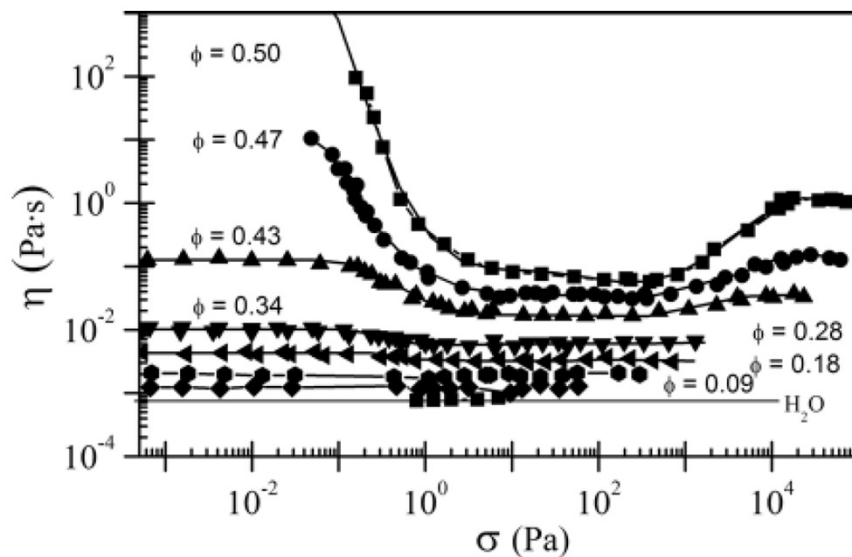


<https://www.youtube.com/watch?v=f2XQ97XHjVw>

https://www.youtube.com/watch?v=hP88C-_LgnE

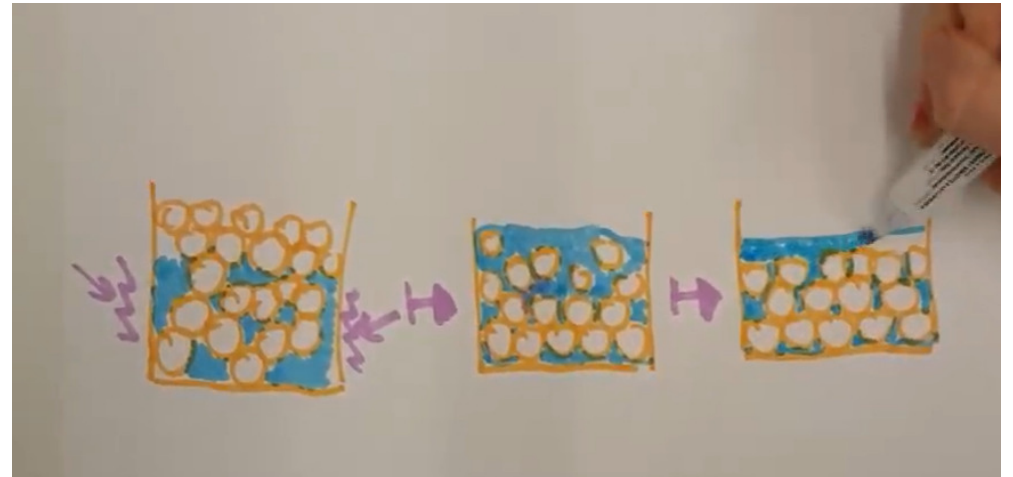
Shear thickening

- Can be observed in suspensions of particles (spherical or nonspherical)
 - Slow deformation: alignment of particles in flow direction (thinning)
 - Fast deformation: jamming and solidification (thickening)



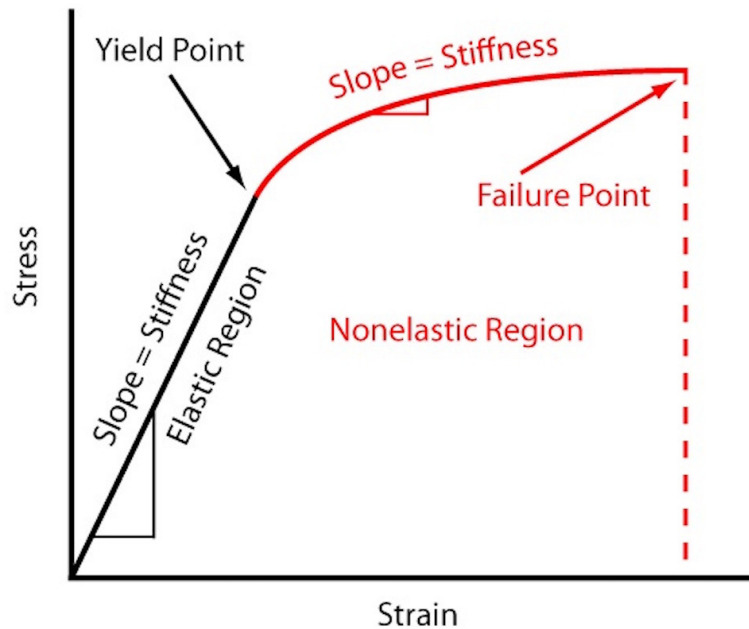
Reynolds dilation

- If the equilibrium state of particle suspension is already solid...
... and you start to shake gently

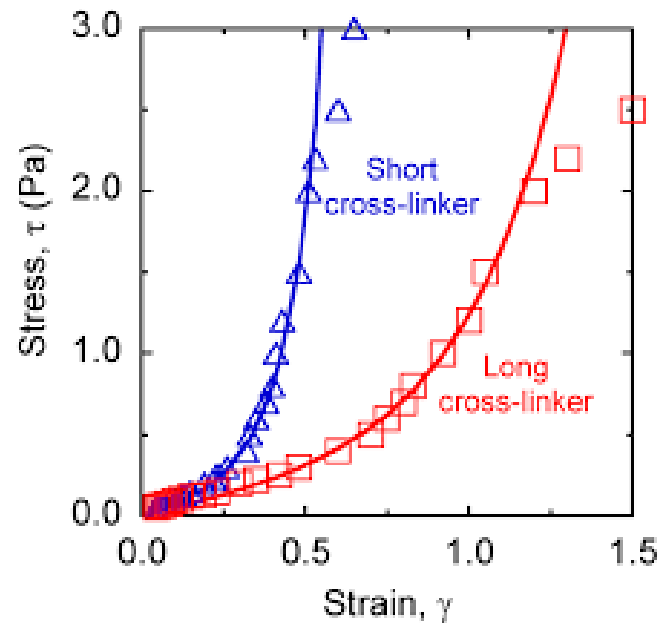


Strain stiffening

- Can be observed in crosslinked networks of rigid polymers



Normal, "yielding" network

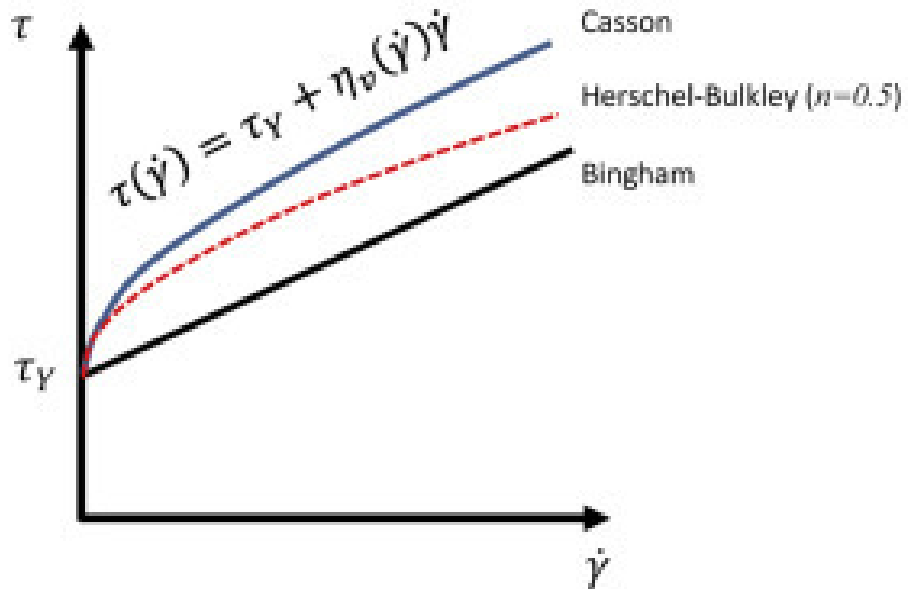


Strain stiffening network



Yield stress

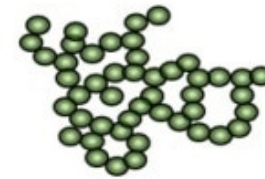
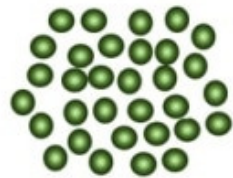
- Common in many microstructured soft materials



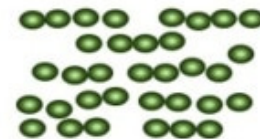
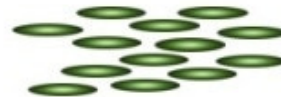
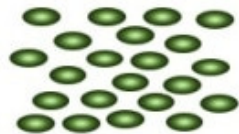
Yield stress

-

microstructure at rest



microstructure under shear



Polymer chains
disentangling and
stretching

Emulsion droplets
reorganising and
deforming

Elongated particles
aligning with the
flow

Aggregated structures
breaking down to
primary particles

Thixotropy and rheopey

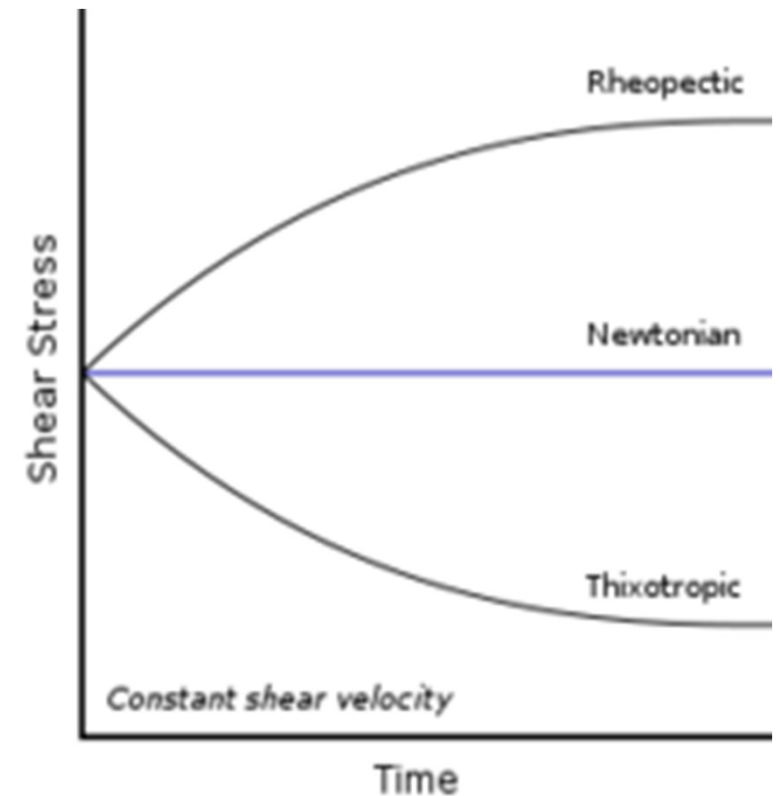
- Time-dependence and hysteresis effects



Ketchup is a good example of a thixotropic substance. It does not flow even when the bottle is held upside down



However, when the bottle is struck (vibrated), the ketchup temporarily liquefies and flows



Thixotropy and rheopecty

- Time-dependence and hysteresis effects

